
Quantum Mechanics

Physics 237

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Solving the Schrödinger Equation using Mathematica.

In[1]:= Solve the finite potential well problem using Mathematica.
The solution is based on the following text book: **Mathematica for Theoretical Physics, Gerd Baumann.**

In[1]:= (* Solve the Schroedinger equation in regions 1 and 3 $x < -a$ and $x > a$. *)
(* The solution is stored in s13. The constant κ is defined as $\kappa = \sqrt{2m(V_0 - E)}/\hbar$ *)
s13 = DSolve[$\partial_{x,x} \psi[x] - \kappa^2 \psi[x] = 0$, ψ , x] // Flatten

(* Solve the Schroedinger equation in region 2: $-a < x < a$. *)
(* The solution is stored in s2. The constant k is defined as $k = \sqrt{2mE}/\hbar$ *)
s2 = DSolve[$\partial_{x,x} \psi[x] + k^2 \psi[x] = 0$, ψ , x] // Flatten

Out[1]:= $\{\psi \rightarrow \text{Function}[\{x\}, e^{x \kappa} c_1 + e^{-x \kappa} c_2]\}$

Out[2]:= $\{\psi \rightarrow \text{Function}[\{x\}, c_1 \text{Cos}[k x] + c_2 \text{Sin}[k x]]\}$

In[3]:= (* Construct the wavefunctions and define the appropriate integration constants in the various regions. *)

(* In region 1, we must ensure that the wavefunction remains finite for all x and as a result $C[2]$ must be zero. *)

ps1 = $\psi[x]$ /. s13 /. {C[1] \rightarrow A1, C[2] \rightarrow B1} /. B1 \rightarrow 0

(* In region 3, we must ensure that the wavefunction remains finite for all x and as a result $C[1]$ must be zero. *)

ps3 = $\psi[x]$ /. s13 /. {C[1] \rightarrow A3, C[2] \rightarrow B3} /. A3 \rightarrow 0

(* In region 2, there are no constraints on the integration constants. *)

ps2 = $\psi[x]$ /. s2 /. {C[1] \rightarrow A2, C[2] \rightarrow B2}

In[6]:= (* A valid solution to the Schroedinger equation requires that the wavefunctions and their first derivatives match at $x = -a$ and at $x = a$. *)

(* This matching condition will generate 4 equations that must be satisfied: eq1, eq2, eq3, and eq4. *)

(* Equation 1: match wavefunctions at $x = -a$. *)
eq1 = ps1 == ps2 /. x \rightarrow -a

(* Equation 2: match derivative of wavefunctions at $x = -a$. *)
eq2 = ∂_x ps1 == ∂_x ps2 /. x \rightarrow -a

(* Equation 3: match wavefunctions at $x = a$. *)
eq3 = ps3 == ps2 /. x \rightarrow a

(* Equation 4: match derivative of wavefunctions at $x = a$. *)
eq4 = ∂_x ps3 == ∂_x ps2 /. x \rightarrow a

Out[6]:= $A1 e^{-a \kappa} = A2 \text{Cos}[a k] - B2 \text{Sin}[a k]$

Out[7]:= $A1 e^{-a \kappa} \kappa = B2 k \text{Cos}[a k] + A2 k \text{Sin}[a k]$

Out[8]:= $B3 e^{-a \kappa} = A2 \text{Cos}[a k] + B2 \text{Sin}[a k]$

Out[9]:= $-B3 e^{-a \kappa} \kappa = B2 k \text{Cos}[a k] - A2 k \text{Sin}[a k]$

In[10]:= (* These four equations can be combined and written in the following way:

$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \begin{pmatrix} A1 \\ A2 \\ B2 \\ B3 \end{pmatrix} = 0$$

Solving the Schrödinger Equation using Mathematica.

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There is a non-trivial solution to this equation if the determinant of the
matrix is equal to 0.
*)

(* Construct the determinant of the matrix. *)
det1 =
Map[Coefficient[{eq1, eq2, eq3, eq4} /. Equal[a_, b_] => a - b], #] &, {A1, A2, B2, B3} //
Transpose // Det // Simplify

(* Determine when the determinant is equal to 0.
The solutions are stored in the spectral array. *)
spectral = MapAll[PowerExpand[#] &, Simplify[Flatten[Solve[det1 == 0, κ]]] // FullSimplify

(* Determine the coefficients of the wavefunction for the first solution. *)
sol1 = Solve[{eq1, eq2, eq3, eq4} /. spectral[[1]], {A1, B2, A2, B3}] // Simplify // Flatten

(* Determine the coefficients of the wavefunction for the second solution. *)
sol2 = Solve[{eq1, eq2, eq3, eq4} /. spectral[[2]], {A1, A2, B2, B3}] // Simplify // Flatten

Out[10]=  $e^{-2 a \kappa} (2 \kappa \cos[2 a \kappa] + (-\kappa^2 + \kappa^2) \sin[2 a \kappa])$ 
Out[11]=  $\{\kappa \rightarrow -\kappa \cot[a \kappa], \kappa \rightarrow \kappa \tan[a \kappa]\}$ 
Out[12]=  $\{B2 \rightarrow -A1 e^{a \kappa \cot[a \kappa]} \csc[a \kappa], A2 \rightarrow 0, B3 \rightarrow -A1\}$ 
Out[13]=  $\{A2 \rightarrow A1 e^{-a \kappa \tan[a \kappa]} \sec[a \kappa], B2 \rightarrow 0, B3 \rightarrow A1\}$ 

In[14]= (* In each region we obtain two solutions:
one symmetric and one asymmetric solution.
*)

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In[14]= (* In each region we obtain two solutions:
one symmetric and one asymmetric solution.
*)

(* Start with region 1: x < -a *)
ψ1a = ps1 /. sol1 /. spectral[[1]]
ψ1s = ps1 /. sol2 /. spectral[[2]]

(* Continue with region 2: -a < x < a *)
ψ2a = ps2 /. sol1 /. spectral[[1]]
ψ2s = ps2 /. sol2 /. spectral[[2]]

(* Finish with region 3: a < x *)
ψ3a = ps3 /. sol1 /. spectral[[1]]
ψ3s = ps3 /. sol2 /. spectral[[2]]

Out[14]=  $A1 e^{-\kappa x \cot[a \kappa]}$ 
Out[15]=  $A1 e^{\kappa x \tan[a \kappa]}$ 
Out[16]=  $-A1 e^{a \kappa \cot[a \kappa]} \csc[a \kappa] \sin[\kappa x]$ 
Out[17]=  $A1 e^{-a \kappa \tan[a \kappa]} \cos[\kappa x] \sec[a \kappa]$ 
Out[18]=  $-A1 e^{\kappa x \cot[a \kappa]}$ 
Out[19]=  $A1 e^{-\kappa x \tan[a \kappa]}$ 

In[20]=

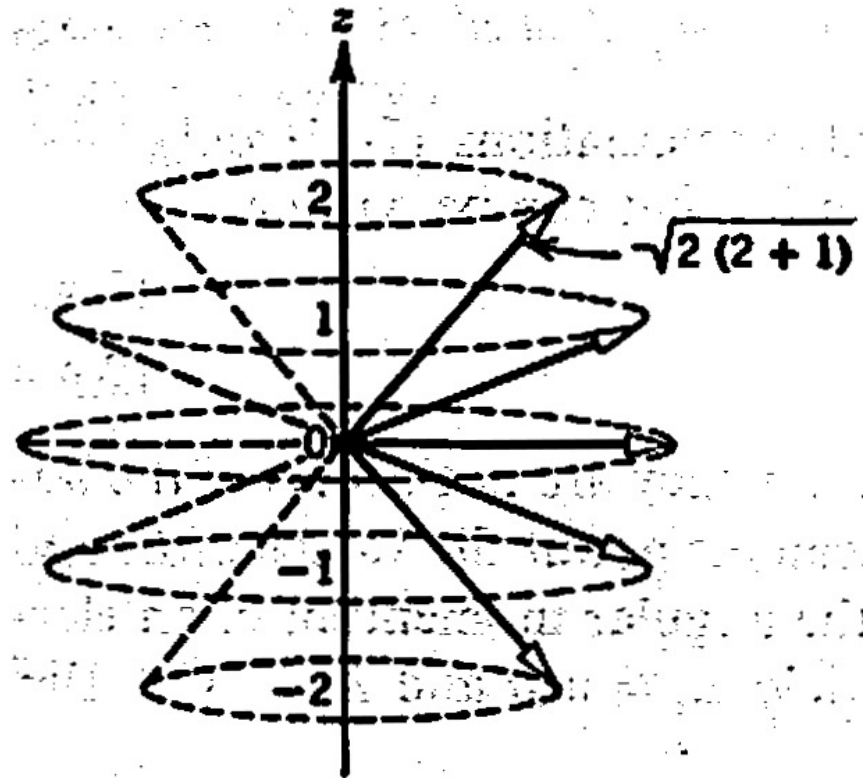
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Screenshot

Final Comments on Spherical Symmetry.

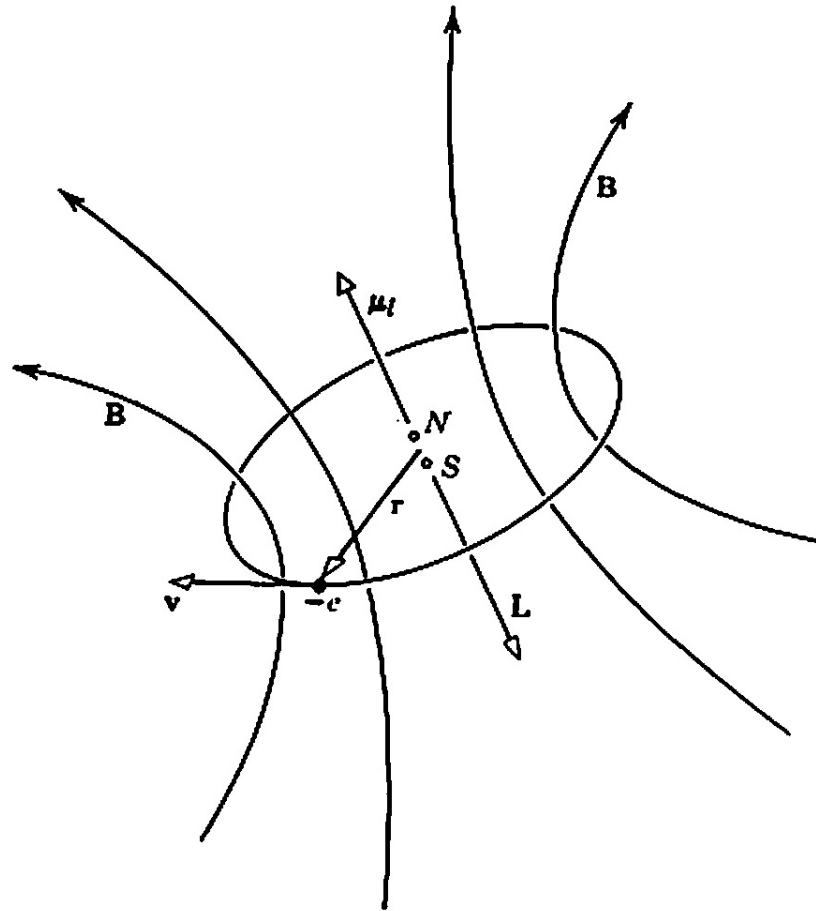
Table 7-2 Some Eigenfunctions for the One-Electron Atom

Quantum Numbers			Eigenfunctions
n	l	m_l	
1	0	0	$\psi_{100} = \frac{1}{\sqrt{\pi}} \left(\frac{Z}{a_0}\right)^{3/2} e^{-Zr/a_0}$
2	0	0	$\psi_{200} = \frac{1}{4\sqrt{2\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \left(2 - \frac{Zr}{a_0}\right) e^{-Zr/2a_0}$
2	1	0	$\psi_{210} = \frac{1}{4\sqrt{2\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \frac{Zr}{a_0} e^{-Zr/2a_0} \cos \theta$
2	1	± 1	$\psi_{21\pm 1} = \frac{1}{8\sqrt{\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \frac{Zr}{a_0} e^{-Zr/2a_0} \sin \theta e^{\pm i\phi}$
3	0	0	$\psi_{300} = \frac{1}{81\sqrt{3\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \left(27 - 18\frac{Zr}{a_0} + 2\frac{Z^2r^2}{a_0^2}\right) e^{-Zr/3a_0}$
3	1	0	$\psi_{310} = \frac{\sqrt{2}}{81\sqrt{\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \left(6 - \frac{Zr}{a_0}\right) \frac{Zr}{a_0} e^{-Zr/3a_0} \cos \theta$
3	1	± 1	$\psi_{31\pm 1} = \frac{1}{81\sqrt{\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \left(6 - \frac{Zr}{a_0}\right) \frac{Zr}{a_0} e^{-Zr/3a_0} \sin \theta e^{\pm i\phi}$
3	2	0	$\psi_{320} = \frac{1}{81\sqrt{6\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \frac{Z^2r^2}{a_0^2} e^{-Zr/3a_0} (3 \cos^2 \theta - 1)$
3	2	± 1	$\psi_{32\pm 1} = \frac{1}{81\sqrt{\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \frac{Z^2r^2}{a_0^2} e^{-Zr/3a_0} \sin \theta \cos \theta e^{\pm i\phi}$
3	2	± 2	$\psi_{32\pm 2} = \frac{1}{162\sqrt{\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \frac{Z^2r^2}{a_0^2} e^{-Zr/3a_0} \sin^2 \theta e^{\pm 2i\phi}$

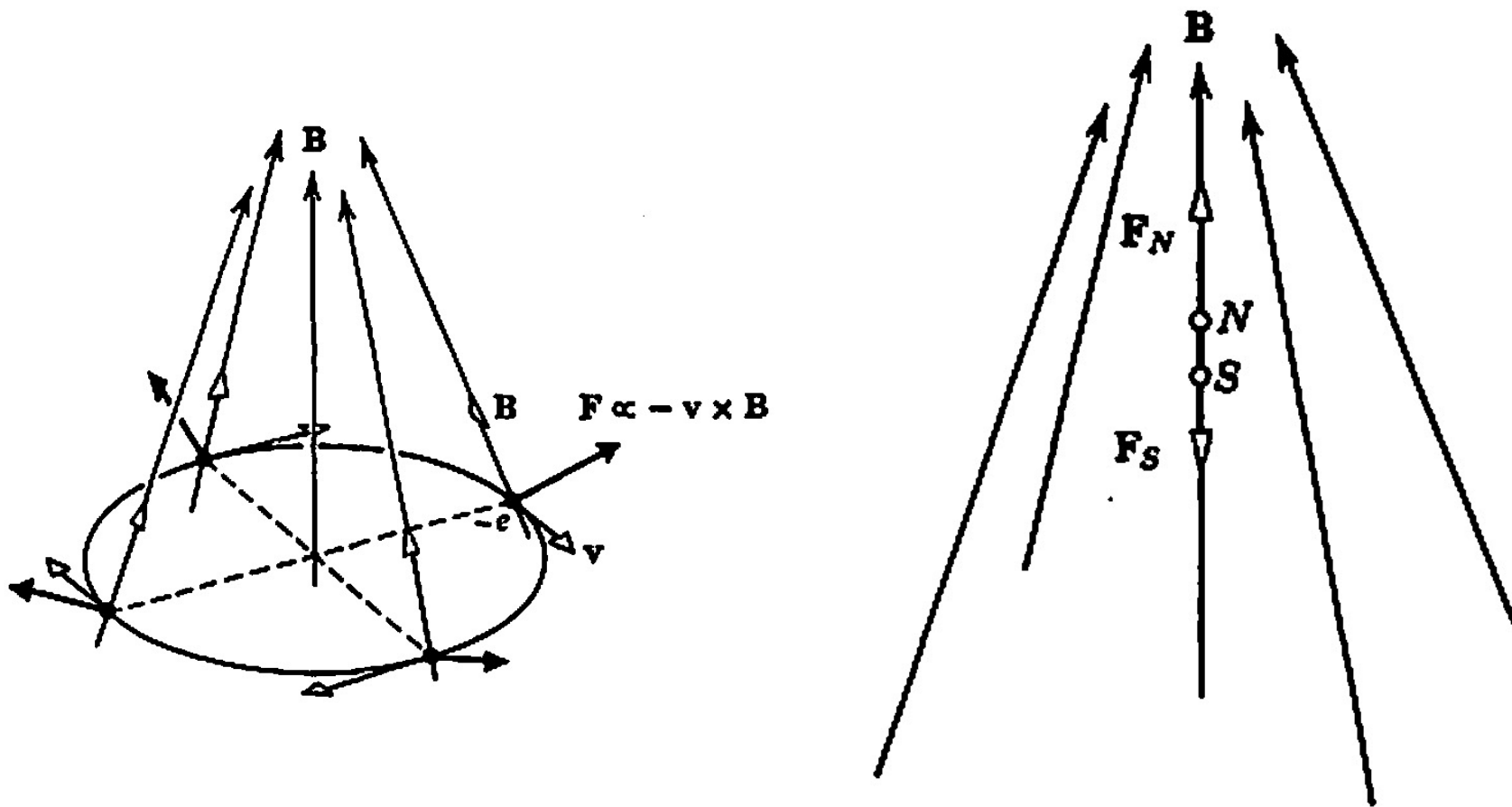


Chapter 8.

An atomic electron in a magnetic field.



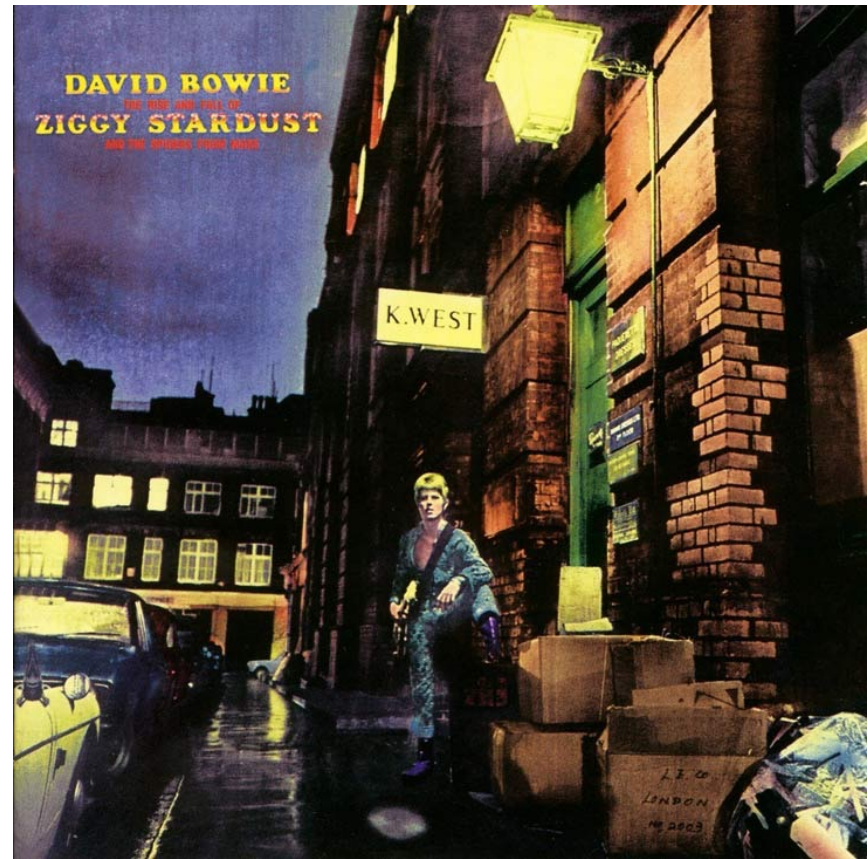
Forces on a magnetic dipole.



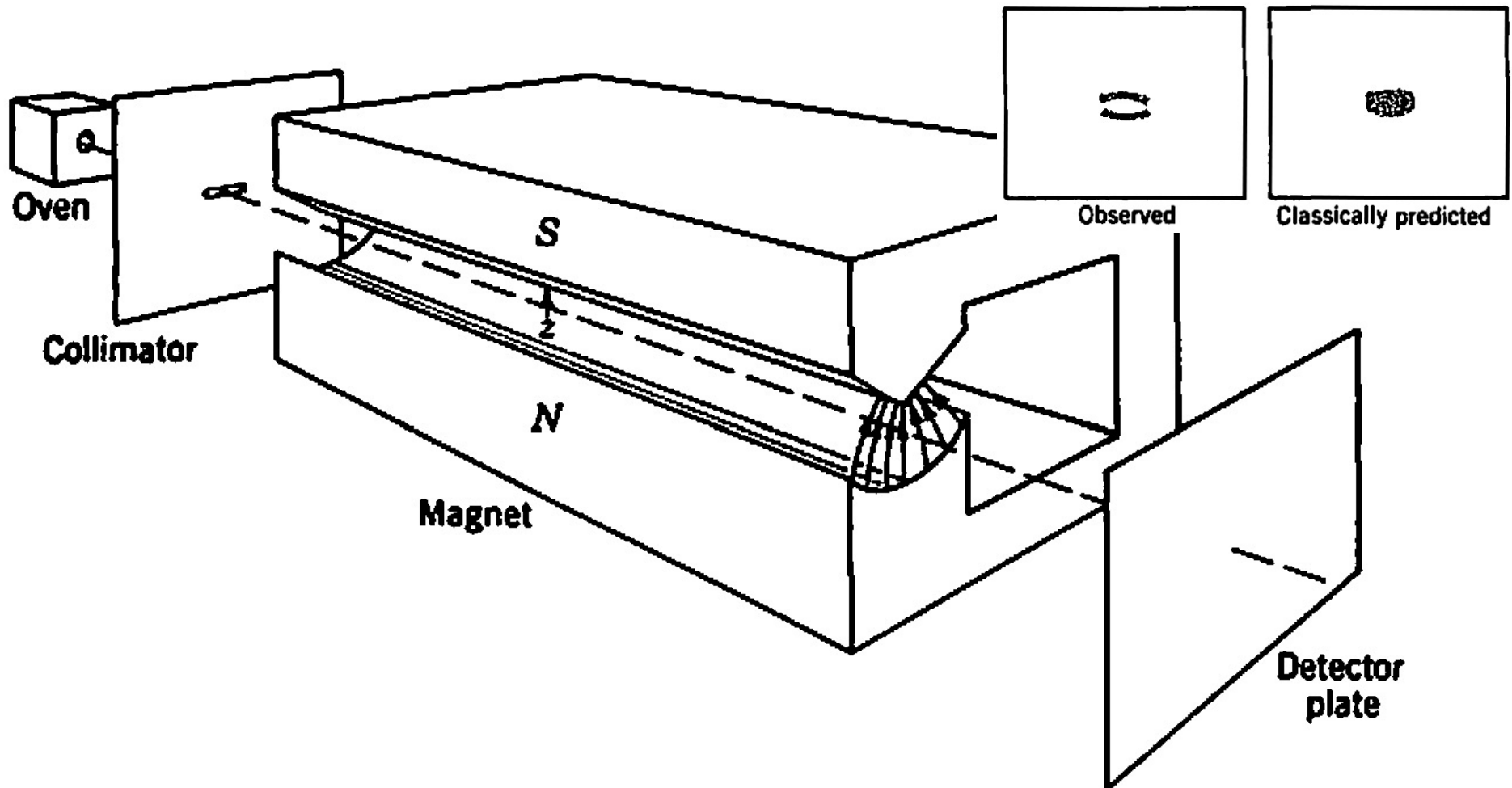


4 Minute 13 Second Intermission.

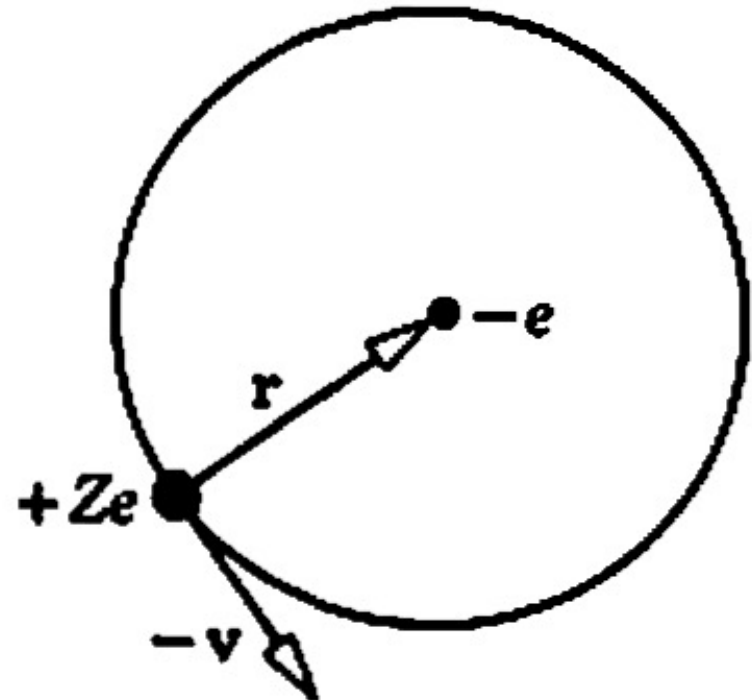
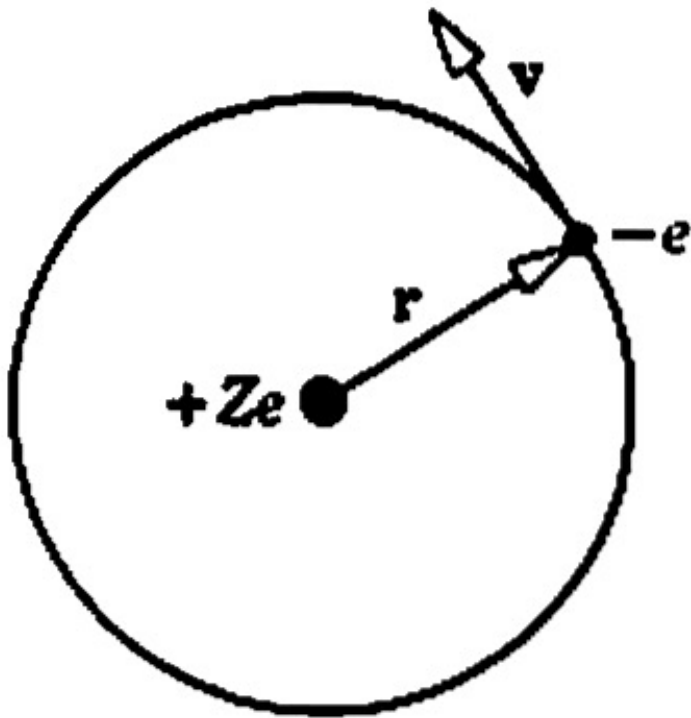
- Since paying attention for 1 hour and 15 minutes is hard when the topic is physics, let's take a 4 minute 13 second intermission.
- You can:
 - Stretch out.
 - Talk to your neighbors.
 - Ask me a quick question.
 - Enjoy the fantastic music.



The Stern-Gerlach Experiment.



The Spin-Orbit Interaction.



ENOUGH FOR TODAY?