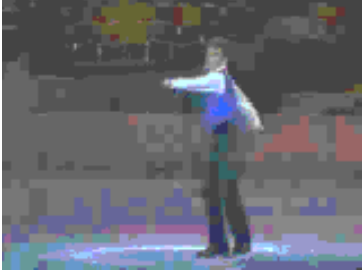


Physics 141.  
Lecture 17.



Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 1

---

---

---

---

---

---

---

---

Physics 141.  
Lecture 17.

- Course information.
- Quiz
- Topics to be discussed today (Chapter 11):
  - Rotational Variables
  - Rotational Kinetic Energy
  - Torque

Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 2

---

---

---

---

---

---

---

---

Physics 141.  
Course information.

- Homework set # 7 is due on Friday 11/4 at noon.
- Homework set # 8 is due on Friday 11/11 at noon.
- Lab report # 4 is due on Wednesday 11/9 at noon.

Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 3

---

---

---

---

---

---

---

---

## The Personal Response System (PRS). Quiz.



Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 4

---

---

---

---

---

---

---

---

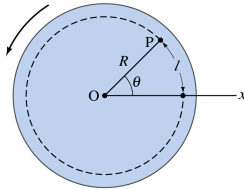
## Rotational variables.

The variables that are used to describe rotational motion are:

- Angular position  $\theta$
- Angular velocity  $\omega = d\theta/dt$
- Angular acceleration  $\alpha = d\omega/dt$

The rotational variables are related to the linear variables:

- Linear position  $l = R\theta$
- Linear velocity  $v = R\omega$
- Linear acceleration  $a = R\alpha$



Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 5

---

---

---

---

---

---

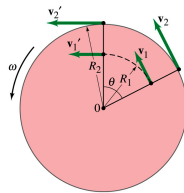
---

---

## Rotational variables.

Things to consider when looking at the rotation of rigid objects around a fixed axis:

- Each part of the rigid object has the same angular velocity.
- Only those parts that are located at the same distance from the rotation axis have the same linear velocity.
- The linear velocity of parts of the rigid object increases with increasing distance from the rotation axis.



Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 6

---

---

---

---

---

---

---

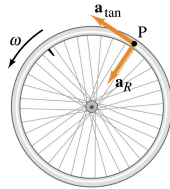
---

## Rotational variables.

• Note: the acceleration  $a_t = r\alpha$  is only one of the two components of the acceleration of point P. The two components of the acceleration of point P are:

• The **radial component**: this component is always present since point P carried out circular motion around the axis of rotation.

• The **tangential component**: this component is present only when the angular acceleration is not equal to  $0 \text{ rad/s}^2$ .



Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 7

---

---

---

---

---

---

---

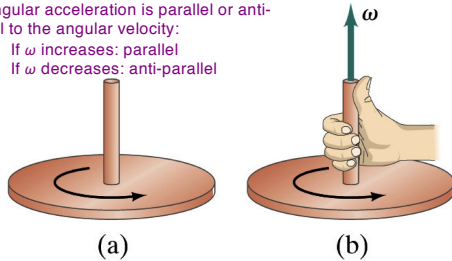
---

## Rotational variables.

Angular velocity and acceleration are vectors! They have a magnitude and a direction. The direction of  $\omega$  is found using the right-hand rule.

The angular acceleration is parallel or anti-parallel to the angular velocity:

If  $\omega$  increases: parallel  
If  $\omega$  decreases: anti-parallel



Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 8

---

---

---

---

---

---

---

---

## Rotational kinetic energy.

• Since the components of a rotating object have a non-zero (linear) velocity we can associate a kinetic energy with the rotational motion:

$$K = \sum_i \frac{1}{2} m_i v_i^2 = \frac{1}{2} \sum_i m_i (\omega r_i)^2 = \frac{1}{2} \left( \sum_i m_i r_i^2 \right) \omega^2 = \frac{1}{2} I \omega^2$$

• The kinetic energy is proportional to the rotational velocity  $\omega$ . Note: the equation is similar to the translational kinetic energy ( $\frac{1}{2} m v^2$ ) except that instead of being proportional to the mass  $m$  of the object, the rotational kinetic energy is proportional to the **moment of inertia  $I$**  of the object:

$$I = \sum_i m_i r_i^2 \quad \text{Note: units of } I: \text{ kg m}^2$$

Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 9

---

---

---

---

---

---

---

---

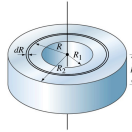
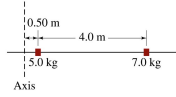
## The moment of inertia $I$ . Calculating $I$ .

- The moment of inertia of an object depends on the mass distribution of object and on the location of the rotation axis.
- For discrete mass distribution it can be calculated as follows:

$$I = \sum_i m_i r_i^2$$

- For continuous mass distributions we need to integrate over the mass distribution:

$$I = \int r^2 dm$$




---

---

---

---

---

---

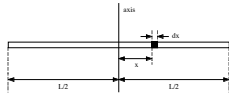
---

---

## Calculating the moment of inertia. Sample problem.

- Consider a rod of length  $L$  and mass  $m$ . What is the moment of inertia with respect to an axis through its center of mass?
- Consider a slice of the rod, with width  $dx$ , located a distance  $x$  from the rotation axis. The mass  $dm$  of this slice is equal to

$$dm = \frac{m}{L} dx$$




---

---

---

---

---

---

---

---

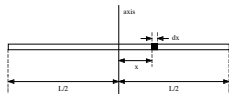
## Calculating the moment of inertia. Sample problem.

- The moment of inertia  $dI$  of this slice is equal to

$$dI = x^2 dm = \frac{m}{L} x^2 dx$$

- The moment of inertia of the rod can be found by adding the contributions of all of the slices that make up the rod:

$$I = \int_{-L/2}^{L/2} \frac{m}{L} x^2 dx = \frac{m}{3L} \left[ \left(\frac{L}{2}\right)^3 - \left(-\frac{L}{2}\right)^3 \right] = \frac{1}{12} mL^2$$




---

---

---

---

---

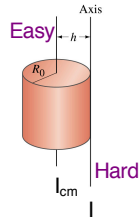
---

---

---

## Calculating the moment of inertia. Parallel-axis theorem.

- Calculating the moment of inertia with respect to a symmetry axis of the object is in general easy.
- It is much harder to calculate the moment of inertia with respect to an axis that is not a symmetry axis.
- However, we can make a hard problem easier by using the parallel-axis theorem:



$$I = I_{cm} + Mh^2$$

Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 13

---

---

---

---

---

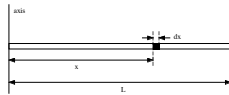
---

---

---

## Calculating the moment of inertia. Sample problem.

- Consider a rod of length  $L$  and mass  $m$ . What is the moment of inertia with respect to an axis through its left corner?
- We have determined the moment of inertia of this rod with respect to an axis through its center of mass. We use the parallel-axis theorem to determine the moment of inertia with respect to the current axis:



$$I = I_{cm} + m\left(\frac{L}{2}\right)^2 = \frac{1}{12}mL^2 + \frac{1}{4}mL^2 = \frac{1}{3}mL^2$$

Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 14

---

---

---

---

---

---

---

---

## 3 Minute 2 Second Intermission.



- Since paying attention for 1 hour and 15 minutes is hard when the topic is physics, let's take a 3 minute 2 second intermission.

- You can:
  - Stretch out.
  - Talk to your neighbors.
  - Ask me a quick question.
  - Enjoy the fantastic music.
  - Solve a WeBWork problem.



Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 15

---

---

---

---

---

---

---

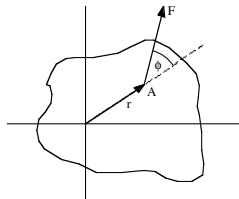
---

## Torque.

- Consider a force  $F$  applied to an object that can only rotate.

- The force  $F$  can be decomposed into two components:

- A **radial component** directed along the direction of the position vector  $r$ . The magnitude of this component is  $F\cos\theta$ . This component will not produce any motion.
- A **tangential component**, perpendicular to the direction of the position vector  $r$ . The magnitude of this component is  $F\sin\theta$ . This component will result in rotational motion.



Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 16

---

---

---

---

---

---

---

---

## Torque.

- If a mass  $m$  is located at the position on which the force is acting (and we assume any other masses can be neglected), it will experience a linear acceleration equal to  $F\sin\theta/m$ .

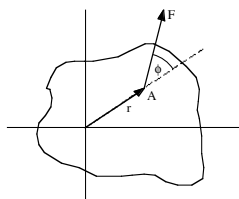
- The corresponding angular acceleration  $\alpha$  is equal to

$$\alpha = \frac{F \sin \theta}{mr}$$

- Since in rotation motion the moment of inertia plays an important role, we will rewrite the angular acceleration in terms of the moment of inertia:

$$\alpha = \frac{rF \sin \theta}{mr^2} = \frac{rF \sin \theta}{I}$$

Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 17




---

---

---

---

---

---

---

---

## Torque.

- Consider rewriting the previous equation in the following way:

$$rF\sin\theta = I\alpha$$

- The left-hand-side of this equation is called the torque  $\tau$  of the force  $F$ :

$$\tau = I\alpha$$

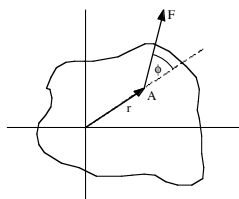
- This equation looks similar to Newton's second law for linear motion:

$$F = ma$$

- Note:

<u>linear</u>	<u>rotational</u>
mass $m$	moment $I$
force $F$	torque $\tau$

Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 18




---

---

---

---

---

---

---

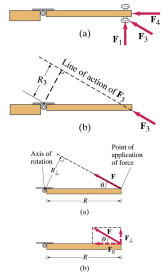
---

## Torque.

- In general the torque associated with a force  $F$  is equal to

$$|\vec{\tau}| = rF \sin\theta = |\vec{r} \times \vec{F}|$$

- The arm of the force (also called the moment arm) is defined as  $r \sin\theta$ . The arm of the force is the perpendicular distance of the axis from the line of action of the force.
- If the arm of the force is 0, the torque is 0, and there will be no rotation.
- The maximum torque is achieved when the angle  $\theta$  is  $90^\circ$ .



Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 19

---

---

---

---

---

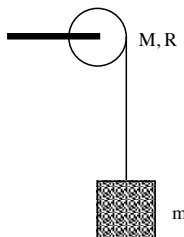
---

---

---

## Rotational motion. Sample problem.

- Consider a uniform disk with mass  $M$  and radius  $R$ . The disk is mounted on a fixed axle. A block with mass  $m$  hangs from a light cord that is wrapped around the rim of the disk. Find the acceleration of the falling block, the angular acceleration of the disk, and the tension of the cord.
- Expectations:
  - Linear acceleration should approach  $g$  when  $M$  approaches 0 kg.



Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 20

---

---

---

---

---

---

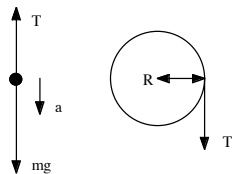
---

---

## Rotational motion. Sample problem.

- Start with considering the forces and torques involved.
- Define the sign convention to be used.
- The block will move down and we choose the positive  $y$  axis in the direction of the linear acceleration.
- The net force on mass  $m$  is equal to

$$ma = mg - T$$



Frank L. H. Wolfs Department of Physics and Astronomy, University of Rochester, Lecture 17, Page 21

---

---

---

---

---

---

---

---

## Rotational motion. Sample problem.

- The net torque on the pulley is equal to

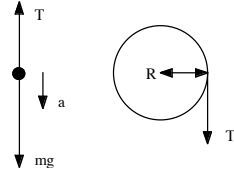
$$\tau = RT$$

- The resulting angular acceleration is equal to

$$\alpha = \frac{\tau}{I} = \frac{RT}{\frac{1}{2}MR^2} = \frac{2T}{MR}$$

- Assuming the cord is not slipping we can determine the linear acceleration:

$$a = \alpha R = 2 \frac{T}{M}$$




---

---

---

---

---

---

---

---

## Rotational motion. Sample problem.

- We now have two expressions for

a:

$$a = 2 \frac{T}{M}$$

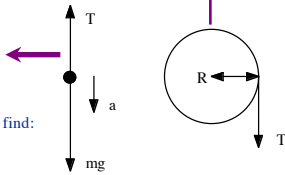
$$a = \frac{mg - T}{m} = g - \frac{T}{m}$$

- Solving these equations we find:

$$T = \frac{M}{M + 2m} mg$$

$$a = \frac{2m}{M + 2m} g$$

**Note:  $a = g$  when  $M = 0$  kg!!!**




---

---

---

---

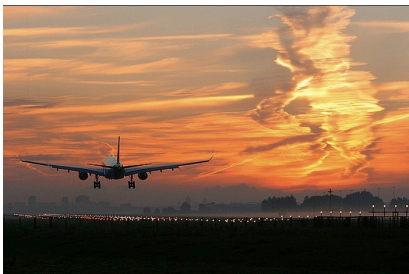
---

---

---

---

## Done for today!



Landing at Amsterdam Airport.

---

---

---

---

---

---

---

---