

## Models of the atom.

## Thomson model:

sphere of positive charge with electrons embedded

- Features:
- 1). Ground state has electrons fixed at equilibrium positions
  - 2). excited states have electrons vibrating around equilibrium positions
  - 3). electron vibration  $\Rightarrow$  emission

note:  $\left\{ \begin{array}{l} \text{qualitative ok} \\ \text{quantitative not ok} \end{array} \right. \Rightarrow \text{kills the model.}$

## Rutherford model:

based on  $\alpha$  scattering of ~~alpha~~ atoms

observed large deflections with finite probability

Facts: since  $m_e \ll m_\alpha$  deflections due to scattering of electron is small

Large deflections require strong

Coulomb repulsion. Can not be

provided if + charge is distributed

over  $10^{-10}$  m.

If large deflections are due to multiple scatterings we can determine the fraction of backwards going  $\alpha$  particles:  $10^{-3500}$

Observed value is  $10^{-4}$

$\Rightarrow$  Thomson + multiple scatterings fails.

$\Rightarrow$  Assume  $\alpha$ 's backscattered due to a single scatter  $\Rightarrow$  require + charge to be concentrated in small volume ( $10^{-14}$  m).

Rutherford developed a model to describe the observed scattering distributions.

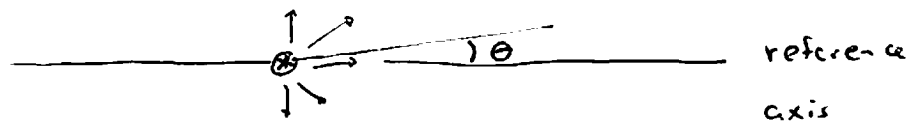
- 1). incident  $\alpha$  particles have a specific impact parameter. (b).
- 2). there is a one-to-one correlation between impact parameter and
  - ⊗ scattering angle  $\Theta$
  - ⊗ distance of closest approach  $R$ .
- 3). smallest value of  $R$  is a head-on collision. ( $R_{\min} = D$ ).

Result:

$$\frac{N(\theta) d\theta}{I} = \frac{\pi}{8} \rho t D^2 \frac{\sin\theta d\theta}{\sin^4(\theta/2)}$$

$\rho$  → density of nuclei in #/cm<sup>3</sup>  
 $t$  → Thickness foil (cm)  
 $D$  →  $D = \frac{1}{4\pi\epsilon_0} \frac{zZe^2}{(\frac{1}{2}Mv^2)}$  = dist. of closest approach in head-on collision.

- Experiment:
- ① fixed cross section of detector will capture different fractions of particles at different angles
  - ② measured rates must be normalized for this type of variations
  - ③ Consider random emission:



- # of particles emitted with  $\theta = 0^\circ$  will be very small (same for  $180^\circ$ )
  - # of particles emitted with  $\theta = 90^\circ$  will be very large
- ⇒ angle dependent probability.

⊗ remove trivial  $\theta$  dependence

$$d\Omega = \frac{\text{area}}{r^2} = 2\pi \sin\theta d\theta$$

⊗ measured values are converted

$$\text{to } \frac{d\sigma}{d\Omega} = \frac{dN}{I_n} \cdot \frac{1}{d\Omega} =$$

$$= \frac{N d\theta}{I_n} \cdot \frac{1}{d\Omega} =$$

$$= \left(\frac{1}{4\pi\epsilon_0}\right)^2 \left(\frac{zZe^2}{2Mv^2}\right)^2 \frac{1}{\sin^4(\theta/2)}$$

$$\sqrt{8\left(\frac{1}{2}\right)^2 \cdot 2}$$

⊗ deviation of  $\left(\frac{d\sigma}{d\Omega}\right)_{\text{Ruth}}$  as function

of distance can be used to determine the size of the nucleus.

If  $D < R_{\text{nucleus}} \Rightarrow$  nuclear force becomes important  $\Rightarrow$  no pure Rutherford scattering.

## Atomic Spectra:

- ⊗ Wave lengths emitted are discrete.
  - ⊗ hydrogen spectrum is simple.
- empirical formula:

$$\text{Balmer: } \lambda = 3646 \frac{n^2}{n^2 - 4}$$

$$R = \frac{1}{\lambda} = R_H \left( \frac{1}{2^2} - \frac{1}{n^2} \right)$$

↓  
Rydberg's constant  
for H

Similar formula were found for other series.

Bohr atoms:

note: mixture  
of classical and  
non-classical physics.

Bohr made the following assumptions:

- ⊗ electrons move in circular orbits
- ⊗  $L_{orbit} = n\hbar$  (only certain orbits are allowed).
- ⊗ orbits are stable - no radiated energy.
- ⊗ radiation emitted when transitions occur between energy levels.

Classical physics is used to determine energy of electrons :

$$E = - \frac{m_e Z^2 e^4}{(4\pi\epsilon_0)^2 (2k^2)} \frac{1}{n^2} \quad n = 1, 2, 3, \dots$$

Energy  $< 0$  . Unbound electron  
if  $E \geq 0$

Note : predictions agree very well with  
exp. values.

↳ model predicted lines that were not  
observed yet, but later confirmed  
to exist.

Note : precise measurements show that a  
correction must be made for the nuclear  
mass. Instead of using  $m_e$  we  
use the reduced mass  $\mu$  where

$$\mu = \frac{mM}{m+M}$$

Rydberg constant changes by 1 part in 2000  
for Hydrogen (biggest change).

Bohr predictions agree with data within  
3 parts of 100,000. Amazing!

Confirmation of energy states in atoms obtained  
by Franck and Hertz.

- ⊛ accelerate electrons
- ⊛ measure current
- ⊛ observe significant dips:

in Hg first dip at 4.5 V  $\Rightarrow$   
lose a lot of electrons

What happens:

4.5 eV  $e^-$  interacts with  
atom and excites it

- ↳  $e^-$  now has energy  
too low to be collected
- ↳ excited atom emits a  
4.5 eV  $\gamma$