

Aside: white dwarfs, novae, and Type Ia supernovae

Sometimes white dwarfs (WDs) are found in binary systems with ordinary stars, in orbits small enough for hydrogen-rich material to fall from the normal star to the WD.

- ☐ The material on the surface of the WD winds up hot and compressed to high densities.
- ☐ When it gets hot and dense enough it can undergo fusion again, which, lacking a stellar envelope around it, tends to be explosive. Two basic types of these explosions:



Artist's conception of classical nova Z Camelopardalis (GALEX/Caltech/NASA)

 Classical nova: about 10⁴⁵ erg released, stars brighten by a factor of 10⁴-10⁵, WD survives, process can repeat.

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White dwarfs, novae, and Type Ia supernovae (continued)

- Type Ia supernova (SNIa): about 10⁵¹ erg released, stars brighten by factor of about 10¹⁰, **WD destroyed**.
- ☐ SNIa happen when the "match" of fusion in material acquired from the normal star ignites fusion of carbon and oxygen nuclei throughout the entire WD.
- ☐ This in turn can happen when the additional material has pushed the mass of the WD up close to the maximum (Chandrasekhar) mass.
- ☐ Since the total fusion-fuel supply is thus about the same for all of them, the energy released (and luminosity) is too: SNIa are bright "standard bombs."
- ☐ This makes SNIa very useful for measuring distances to galaxies, and we will meet them again in that context.

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Special-relativity interlude: the relativity of mass, the speed-of-light limit, and maximum degeneracy pressure

Why is the speed of light the maximum speed that can be reached by electrons? (Remember, this was not part of the two original axioms from which Einstein started...)

 \square Because of the **relativity of mass**: if a body with rest mass m_0 moves at speed V with respect to an observer, the observer will measure a mass

$$m = \frac{m_0}{\sqrt{1 - V^2/c^2}}$$

for the body. (Another result of the Special Theory.)

□ Note similarity to the formula for time dilation: in particular, that the denominator approaches 0, and thus *m* approaches infinity, if *V* approaches *c*.

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D	PRS interlu	de: relativ	ity of mass		
What is the mass of an object that has rest mass 1 gram, and flies by at 0.99c?					
A. 0.14 gram E. 14.2 gram	B. 0.53 gram	C. 1 gram	D. 7.1 gram		

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PRS interlude: relativity of mass

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Take a guess: how fast would an object have to fly for its mass to be a factor of 2 larger than its rest mass? $_$

A.
$$V = \frac{1}{2}c$$
 B. $V = \frac{\sqrt{3}}{2}c$ C. $V = \frac{\sqrt{5}}{6}c$ D. $V = \frac{99}{100}c$ E. $V = 2c$

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The relativity of mass, the speed-of-light limit, and maximum degeneracy pressure (continued)

So, suppose you have a body moving at very nearly the speed of light, and you want it to exceed the speed of light. What can you do?

- ☐ It needs to **accelerate** for its speed to increase.
- $\hfill \Box$ You need to exert a force on it in order to make it accelerate.
- ☐ The force required is, essentially, proportional to the product of mass and acceleration in your reference frame. (Nonrelativistic version of this statement is Newton's second law: force = mass times acceleration.)
- □ But the mass approaches infinity as *V* approaches *c*, and thus an infinite force is required to accelerate it further. There's no such thing as an infinite force (or indeed an infinite anything), so *c* is the ultimate speed limit.

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Final	ımp.	losion	ot	massive	burned-out	stars

Electron degeneracy pressure can hold up a star of mass $1.4 M_{\odot}$ or less against its weight, and do so indefinitely. Stellar cores in this mass range at death become white dwarfs. For heavier stars – $\geq 10 M_{\odot}$ – gravity overwhelms electron degeneracy pressure, and the collapse doesn't stop with the star at planet size.

- □ As the star is crushed past a circumference of 10^4 cm or so, all the electrons and protons in the star are squeezed together so closely that they rapidly combine to form neutrons: $p+e^-$ + energy $\rightarrow n+v_e$.
- ☐ Eventually, then, the collapse might be stopped by the onset of neutron degeneracy pressure.
- ☐ A star whose weight is held up by neutron degeneracy pressure is called a **neutron star**.

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Oppenheimer's theory of neutron stars

Neutron stars were first proposed to exist, and to cause supernovae by their formation, by Zwicky and Baade (1934).

- ☐ First calculations of their sizes: Landau (1938).
- ☐ Neutron stars are analogous to white dwarfs, but the calculation of their structure is much more difficult, since the strong nuclear force and general relativity must be taken into account.
 - For white dwarfs, special relativity suffices because the gravity of these stars is not strong enough to make general relativistic effects substantial.
- ☐ They may also be expected to have a maximum mass, as white dwarfs do. So neutron-star formation prevents black hole formation only up to that maximum mass.

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Oppenheimer's theory of neutron stars (cont'd)

□ First calculation of maximum mass: Oppenheimer and Volkoff (1939). They got $0.7M_{\odot}$; more recent calculations, with improvements in the expression of the nuclear forces, give $1.5\text{-}3\ M_{\odot}$. (We will use $2M_{\odot}$ in this course.)

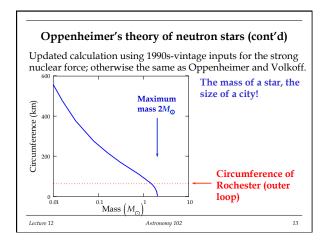


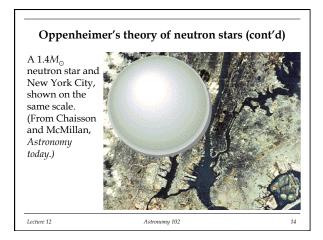
Albert Einstein and J. Robert Oppenheimer at Caltech in 1939. They probably were, at that moment, discussing the prevention of black holes by neutron-star formation.

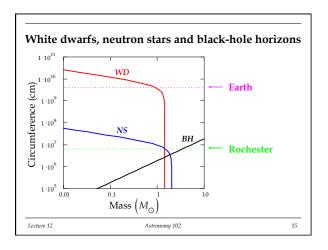
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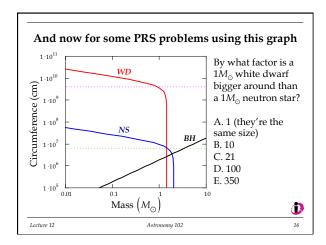
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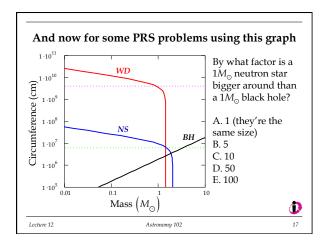
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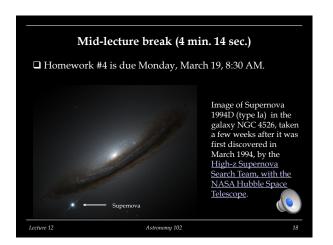












Neutron star formation and Type II supernovae

- ☐ After electron degeneracy pressure is overpowered, and the electrons and protons combine to form neutrons, the star is free to collapse under its weight. Nothing can slow down this collapse until the neutrons are close enough together for their degeneracy pressure to become large.
 - Recall that this requires confinement of each particle to a space a factor of about 1836 smaller than for electron degeneracy pressure.
- ☐ This collapse takes very little time, and the collapsing material is moving very fast when neutron degeneracy pressure hits.
- ☐ A **neutron-degeneracy-pressure supported core** can form from the inner part of the collapsing material.

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Neutron star formation and Type II supernovae (continued) Interstellar dust grain falling from a large distance to star's surface v = 620 km/sec v = 7400 km/sec v = 120,000 km/sec = 0.4c $\Delta E = 1 \text{ erg}$ $\Delta E = 140 \text{ erg}$ $\Delta E = 44,000 \text{ erg}$

Neutron star formation and Type II supernovae (continued)

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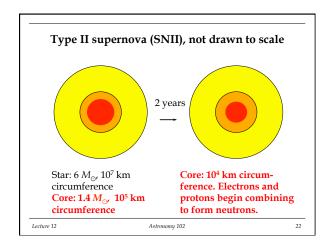
- ☐ The outer, collapsing material that didn't make it into the neutron core proceeds to bounce off this core, rebounding into the rest of the star and exploding it with great violence. This is called a Type II supernova (SNII).
- ☐ The bounce allows much of the energy gained in the collapse to be released, upwards of 10⁵² erg.
- Artist's conception of the formation of a neutron star and a Type II supernova (CXO/CfA/NASA).
- ☐ Because this process will work for just about any star 8 solar masses and heavier, the released energy varies a lot more from SNII to SN II than from SNIa to SNIa.

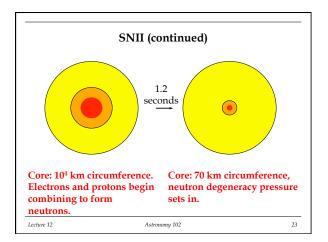
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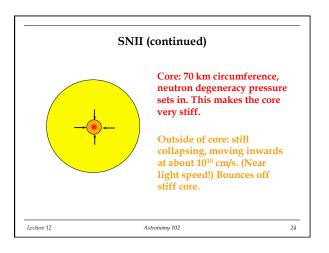
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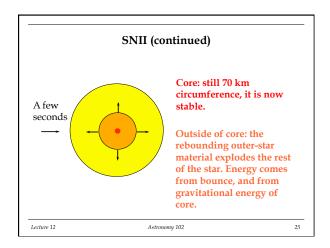
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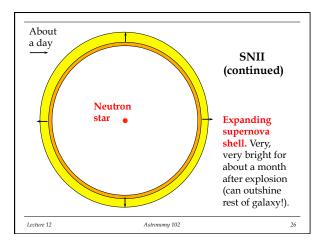
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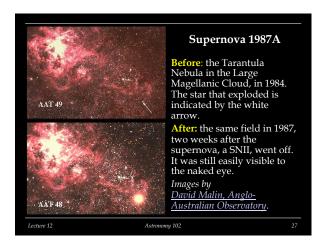


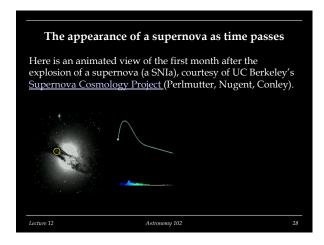












Neutron stars, SNII, and pulsars

Many hundreds of neutron stars are known today; they appear mostly as **pulsars**: starlike sources of radio and visible light whose light output pulsates rapidly.

- ☐ Discovered in 1967 by Cambridge grad student Jocelyn Bell, they were almost immediately identified as rapidlyrotating neutron stars that emit "beams" of light by accelerating electrons and protons outward along their magnetic poles. (They pulse like a lighthouse does.)
- ☐ Many young supernova remnants contain pulsars.
- \square Several pulsars occur in binary systems, for which masses can be measured accurately; all turn out to be around $1.4\text{-}1.5M_{\odot}$, comfortably less than $2M_{\odot}$, and greater than the maximum white-dwarf mass.

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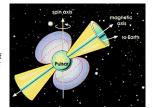
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Neutron stars and pulsars.

- One teaspoonful of a neutron star would weigh a billion tons.
- When a neutron star is formed, it in general spins.
- We can see the neutron stars because most of them have jets of particles bouncing off them, moving away with the speed of light along the magnetic axis of the neutron star.
- Since the star spins, we see this beam of light only when it is directed towards the Earth.
- The rate of the pulsar allows us to measure the spin rate.



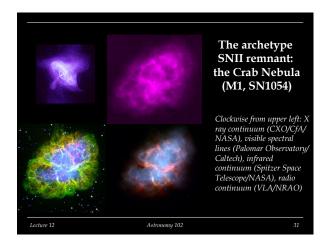
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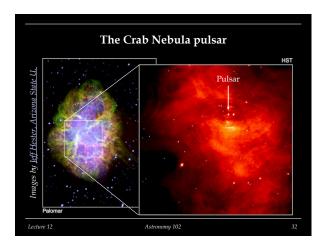
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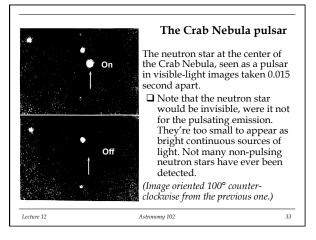
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The Crab Nebula pulsar						
in X-rays, w	(i.e. images taken 0.015 sec apart), but this timith the Chandra X-ray Observatory (CXO), and ientation as page 31.					
	MAIN PULSE "OFF" PHASE PULSAR IN THE CRAB NEBULA					
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Summary: status of the Schwarzschild singularity and black holes

- \square Electron and neutron degeneracy pressure can prevent the formation of black holes from dead stars, but only for core masses below about $2M_{\odot}$.
- ☐ Stars with masses in excess of this must eject material during their final stages of life if they are to become white dwarfs or neutron stars.
- ☐ Judging from the large numbers of white dwarfs, neutron stars, planetary nebulae and supernova we see, the vast majority of stars do end their lives in this way.
- ☐ For core masses larger than this, a pressure stronger than the maximum neutron degeneracy pressure is required to prevent the formation of black holes.

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Summary: status of the Schwarzschild singularity and black holes (continued)

- ☐ There aren't any elementary particles, heavier than neutrons and not radioactive at these high densities, that could provide a larger maximum degeneracy pressure.
- □ In fact, no force known to science exists that would prevent the collapse of a star with a core mass greater than 2M_☉ from proceeding to the formation of a black hole. For very heavy stars, black hole formation is compulsory. The Schwarzschild singularity is real!
- ☐ Einstein was very disappointed in this result, and never trusted it; but had he lived ten years longer, experiments and observational data would have compelled his acceptance.

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